



AS WE EXPLAINED LAST TIME
GEOMETRIC LANGLANDS COMES
FROM MIRROR SYMMETRY AMONG
BRANES

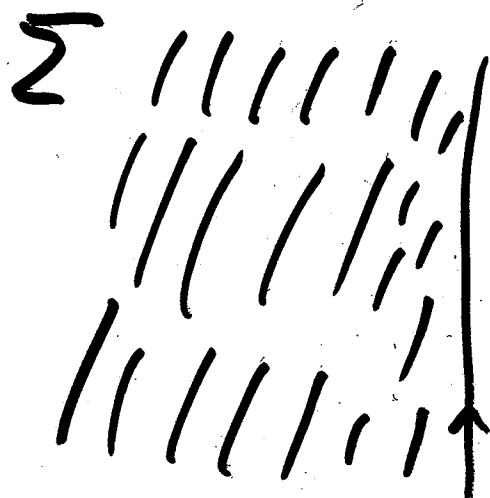
|||X|+|

$M_H(G, C)$
B-BRANE

|||X|||

$M_H(G, C)$
A-BRANE

CONCRETELY THE A-BRANE
OR B-BRANE IS A BOUNDARY
CONDITION IN THE QUANTUM
THEORY



EXAMPLE:
B-BRANE

DIRICHLET BOUNDARY
CONDITIONS

$$\Phi|_{\partial\Sigma} : \partial\Sigma \rightarrow \text{GIVEN POINT IN } \mathcal{M}_H$$

MORE GENERALLY

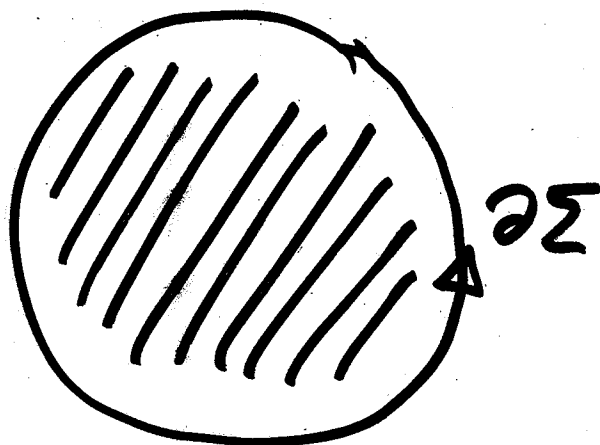
$$\Phi|_{\partial\Sigma} : \partial\Sigma \rightarrow Y \subset M_H$$

Y ENDOWED WITH

VECTOR BUNDLE \mathcal{V}
WITH CONNECTION

PATH INTEGRAL HAS A FACTOR

$$\int \mathcal{D}\Phi \dots \exp(-I) \underbrace{T \exp(-\int_{\partial\Sigma} B)}_{\text{HOLONOMY AROUND } \partial\Sigma}$$



HOLONOMY
AROUND $\partial\Sigma$

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THE SUBMANIFOLD Y AND ITS
VECTOR BUNDLE V OBEY APPROPRIATE
CONDITIONS

B-MODEL THEY DEFINE A
COHERENT SHEAF $\rightarrow X$
M.H.

A-MODEL STANDARD EXAMPLE

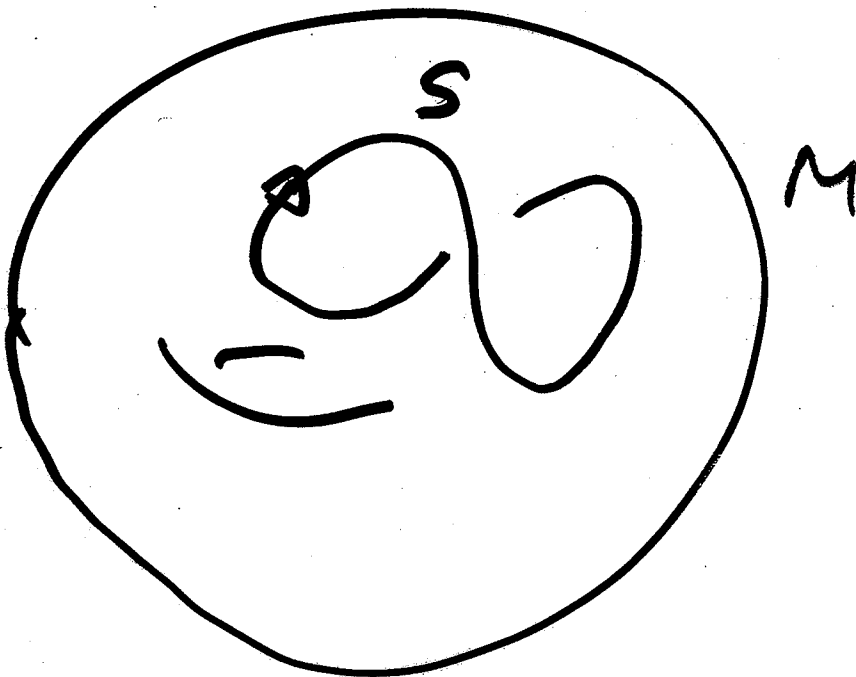
$Y = \text{LAGRANGIAN}$
AND V IS FLAT

PLUS OTHER EXAMPLES
IN GENERAL

("COISOTROPIC BRANES"
- KADUŠIN & PORLOV)

"LINE OPERATORS" IN FOUR 78 DIMENSIONS

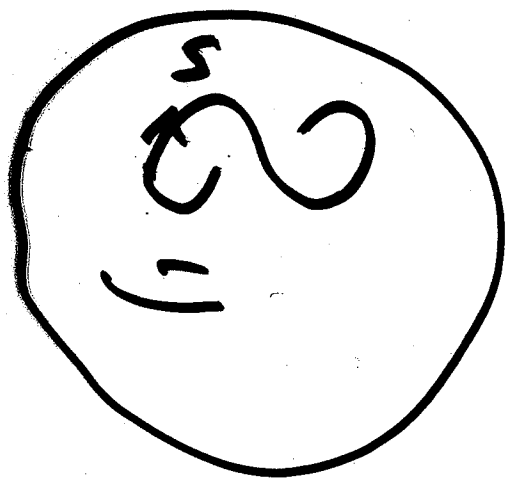
WILSON OPERATOR = HOLONOMY IN
REPRESENTATION R



$$W_R(S) = \text{Tr}_R P \exp \left(- \int_S A \right)$$

= trace of holonomy
in R -rep.

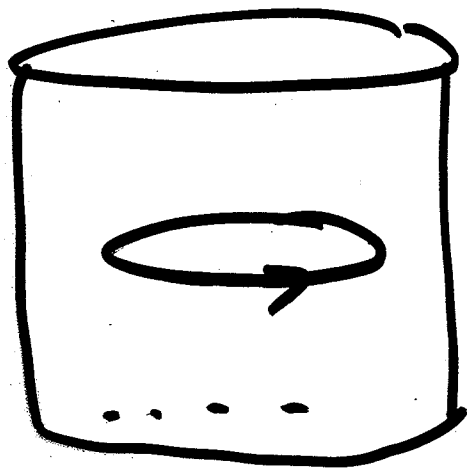
WE INCLUDE THE "WILSON OPERATOR" AS A FACTOR IN PATH INTEGRAL



$$\int_{Q/S} DA d... \exp(-I) \cdot W_R(S)$$

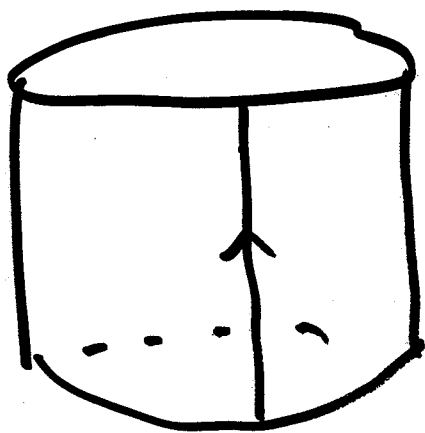
IT IS CALLED AN "OPERATOR" BECAUSE IN A CERTAIN SITUATION IT IS UPON QUANTIZATION A HILBERT SPACE OPERATOR

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Φ TIME

WE'LL ACTUALLY BE INTERESTED
IN THE OPPOSITE SITUATION



DON'T TAKE
TRACE OF
THE HOLONOMY

JUST USE HOLONOMY OPERATOR
FROM INITIAL TO FINAL STATE.

THE "OPERATOR" IS REALLY PART OF
THE DEFINITION OF THE QUANTUM
PROBLEM.

BUT WHAT DOES A "WILSON OPERATOR"⁽⁸¹⁾
TRANSFORM INTO UNDER S-DUALITY?

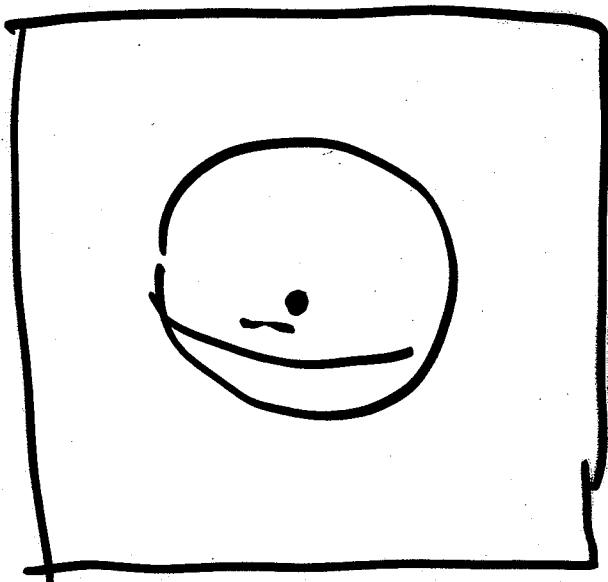
ANSWER (LATE 1970's):

AN "E HOOFT OPERATOR"

TO DEFINE IT, START WITH THE

"MAGNETIC MONOPOLE" SOLUTION

OF MAXWELL'S EQUATIONS ON $\mathbb{R}^3 \setminus 0$



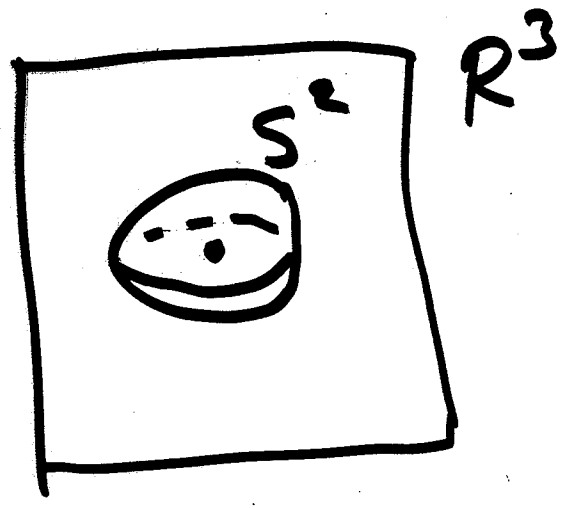
$$F = \frac{1}{2} * d\left(\frac{1}{|x|}\right)$$

S^2

$$\int_{S^2} F = 2\pi$$

THE COEFFICIENT IS CHOSEN

SO THAT, RESTRICTED TO S^2



F IS THE
 CURVATURE OF A
 CONNECTION ON
 $O(1) \rightarrow S^2 \simeq \mathbb{C}P^1$

NOW WE PICK A REPRESENTATION

ρ OF \mathfrak{g} AND A CORRESPONDING

$$\rho: U(1) \rightarrow T \hookrightarrow G$$

THIS GIVES A SOLUTION OF

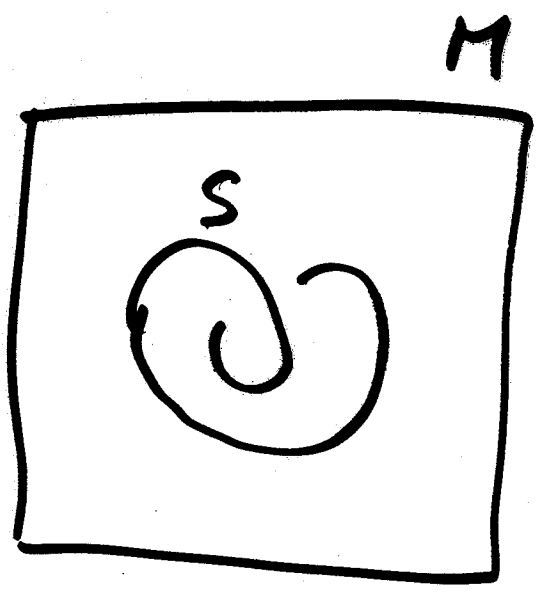
~~MAXWELL~~ YANG-MILLS EQN'S
 FOR G .

NOW DEFINE THE 't HOOFT

OPERATOR $T(S; \mathcal{L}_R)$

IN G GAUGE THEORY

$E \rightarrow M/S$



BY SAYING THAT NEARS THE FIELDS MUST HAVE THE

SINGULARITY DEFINED BY \mathcal{L}_p :

$$\int_{(a_{x...})_{\mathcal{L}_p}} \mathcal{D}A \dots \exp -I$$

THE "USUAL" PATH INTEGRAL, BUT WE INTEGRATE OVER A DIFFERENT SET OF FIELDS

THE S-DUALITY CONJECTURE

SAYS THAT



A WILSON OPERATOR

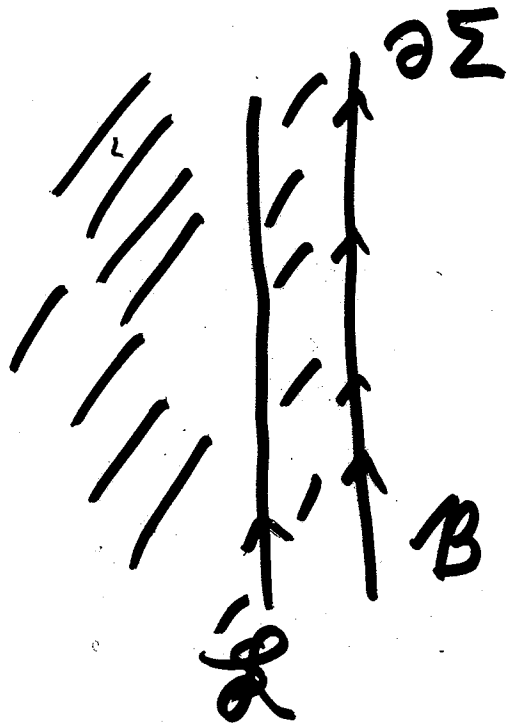
$W(S, L_R)$ IN LG GAUGE

THEORY MAPS TO AN 'T HOOFT

OPERATOR $T(S, L_R)$

IN G GAUGE THEORY.

$\mathcal{L} = W$ OR T



NOW REGARDLESS OF WHETHER
 OUR LINE OPERATOR \mathcal{L} IS A WILSON
 OR 't HOOFT OPERATOR, THINK
 OF IT RUNNING PARALLEL TO
 THE BOUNDARY AS SHOWN.

THE BOUNDARY IS LABELED BY A
 BRANE B .

THE COMPOSITE OF THE LINE
OPERATOR \mathcal{L} AND BRANE \mathcal{B}
IS EQUIVALENT TO A NEW BRANE
THAT WE CALL \mathcal{B}'



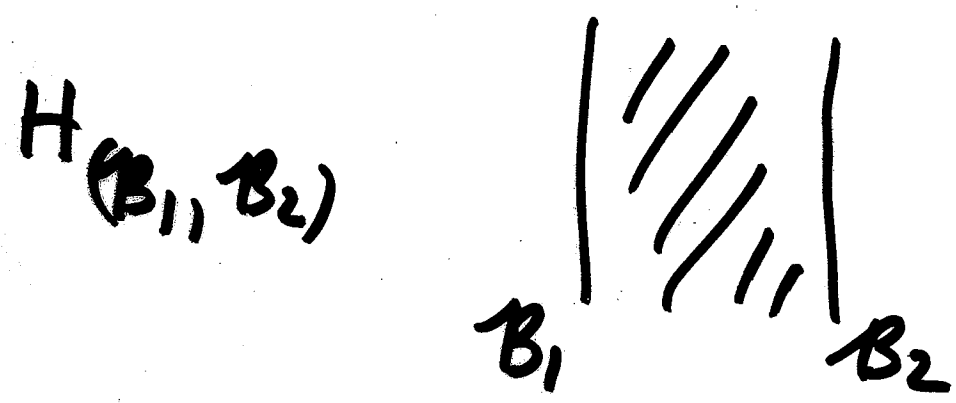
WE WRITE

$$\mathcal{B}' = \mathcal{L} \cdot \mathcal{B}$$

BRANES FORM A "CATEGORY"

GIVEN B_1, B_2 WE DEFINE

A "SPACE OF MORPHISMS"



A LINE OPERATOR \mathcal{L}

GIVES A MORPHISM

FROM THE CATEGORY

OF BRANES TO ITSELF



$$B \rightarrow \mathcal{L} B = B'$$

WE SAY \mathcal{B} IS AN "EIGENBUNDLE" FP
FOR \mathcal{L} IF \mathcal{B}' IS A "MULTIPLE"
OF \mathcal{B} , WHICH MEANS

$$\mathcal{B}' = \mathcal{B} \otimes U$$

U A FIXED VECTOR SPACE

THIS MAKES SENSE ABSTRACTLY

$$\mathcal{H}(\tilde{\mathcal{B}}, \mathcal{B}') = \mathcal{H}(\tilde{\mathcal{B}}, \mathcal{B}) \otimes U$$

FOR ALL $\tilde{\mathcal{B}}$.

BUT IN GEOMETRY IT MEANS

THAT THE SHEAVES \mathcal{V}' , \mathcal{V}

DEFINING \mathcal{B}' , \mathcal{B} OBEY $\mathcal{V}' = \mathcal{V} \otimes U$

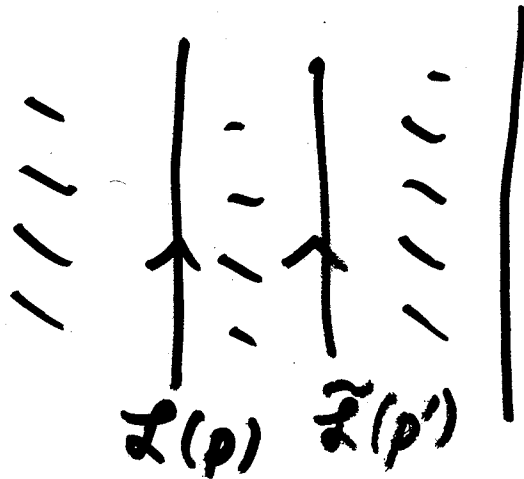
IN TOPOLOGICAL FIELD THEORY

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LINE OPERATORS DERIVED

FROM FOUR DIMENSIONS

COMMUTE WITH EACH OTHER



p, p' TWO POINTS IN C

AND CAN HAVE SIMULTANEOUS

EIGENBRANES

SO AT $\Psi = \infty$, WE

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CAN LOOK FOR A SIMULTANEOUS

EIGENBRANE FOR ALL

WILSON OPERATORS

PEC
$$W(S, \rho, L, R) \mathcal{B} = \mathcal{B} \otimes U_{\rho}(L, R)$$

FOR ALL L, R

WE CALL SUCH A BRANE

AN ELECTRIC EIGENBRANE.

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S-DUALITY WILL MAP AN
"ELECTRIC EIGENBRAVE" TO A
"MAGNETIC EIGENBRAVE":

$$T(S, L_R) \tilde{B} = \tilde{B} \otimes U(L_R)$$

WITH THE SAME

"EIGENVALUE"

THE ZEROBRANE B

DERIVED BY

$\rho: \pi_1(C) \rightarrow LG$

IS AN ELECTRIC EIGENBRANE

SO ITS DUAL

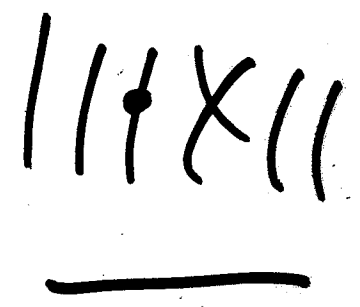
\tilde{B} IS A MAGNETIC

EIGENBRANE, WHICH GEOMETERS

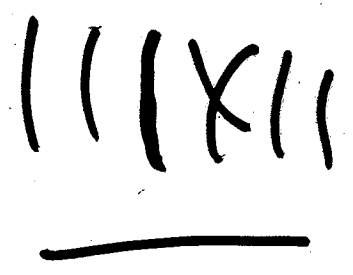
DESCRIBE BY TALKING ABOUT A

"HECKE EIGENSHEAF."

g_{M4}



g_{M4}



TO SHOW THE ZEROBRAKE IS AN
"ELECTRIC EIGENBRAKE": 93
WILSON OPERATOR

$W(p; \rho_R)$

$p = \text{POINT IN } C$

$\rho_R = \text{REP OF } \rho_G$

ACTS AS FOLLOWS:

LET $E_{\rho_R} \rightarrow \mathcal{M}_H(\rho_G) \times C$

BE THE UNIVERSAL BUNDLE

IN REPRESENTATION ρ_R

AND RESTRICT IT TO

~~MANIFOLD~~

$\mathcal{M}_H \times P$

~~CORRESPONDING~~

(94)

$$\mathcal{O}_M \times P \quad P \in C$$

I CALL THE RESTRICTION

$$\mathcal{E}_{LR|P} \rightarrow \mathcal{O}_M$$

THEN IN GENERAL FOR A

BRANE B DERIVED BY A

SHEAF $\mathcal{Q} \rightarrow \mathcal{O}_M(LG, C)$

THE LINE OPERATOR

$W(P; LR)$ ACTS BY

$$\mathcal{Q} \rightarrow \mathcal{Q} \otimes \mathcal{E}_{LR|P}$$

SO B IS AN ELECTRIC
EIGENBRANE PRECISELY IF

$$\mathcal{E}_{LR|P}$$

I.E. THE UNIVERSAL BUNDLE

$$\mathcal{E}_{LR} \rightarrow \mathcal{M}_H \times C$$

RESTRICTED TO $\mathcal{M}_H \times P$

IS TRIVIAL IF FURTHER RESTRICTED
TO THE SUPPORT OF \mathcal{V}

THIS IS CERTAINLY SO IF \mathcal{V}
IS SUPPORTED AT A POINT IN \mathcal{M}_H
CORRESPONDING TO $\rho: \pi_1(C) \rightarrow G_C$

||X||

IN THAT CASE, $W(LR, p)$
ACTS BY

$$\mathcal{V} \rightarrow \mathcal{V} \otimes \mathcal{E}_{LR} / \mathcal{P} \times \mathcal{P}$$

AND THE "EIGENVALUE" IS THE
FIBER AT \mathcal{P} OF THE FLAT

BUNDLE

$$U_{LR} \rightarrow \mathbb{C}$$

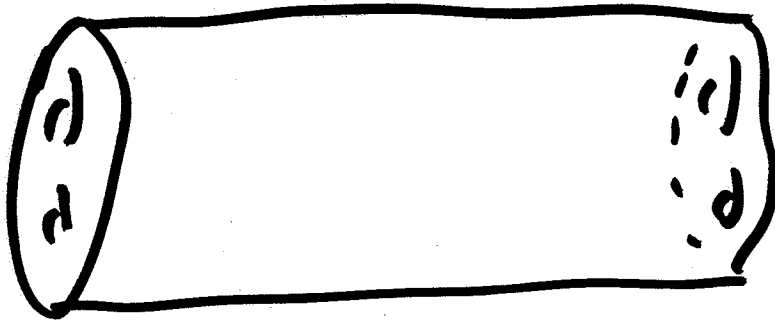
ASSOCIATED TO \mathcal{P} IN THE REPRESENTATION
ACTION
 LR .

SO AFTER DUALITY, THE DUAL
BRANE \tilde{B}

$$\underline{||| \times |||}$$

WILL BE AN EIGENBRANE FOR THE
1/2 HOOFT OPERATORS WITH THE
SAME "EIGENVALUE"

HOW DO 't HOOFT OPERATORS
ACT ON BRAVES?



$$W = I \times C$$

THE DERIVATION LEADS US TO
NONLINEAR PDE'S ASSOCIATED
WITH SUPERSYMMETRY.

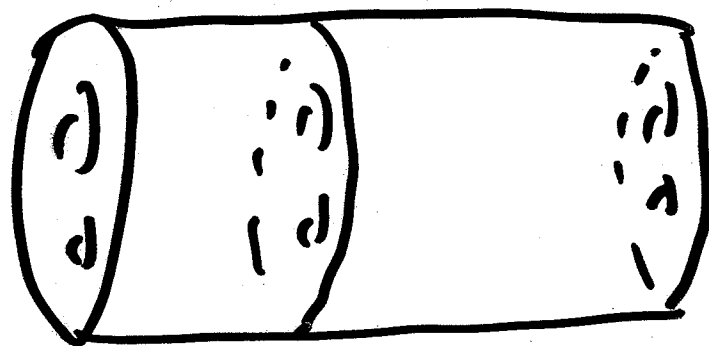
IF WE OMIT THE HIGGS
FIELD, THESE BECOME EQNS
FAMILIAR IN DONALDSON THEORY
AND RELATED SUBJECTS

BOGOMOLNY EQUATIONS IN THREE DIMENSIONS

$$F = * D\phi$$

FOR CONNECTION A ON
 G -BUNDLE $E \rightarrow W =$
 3 -MANIFOLD

AND $\phi \in \Omega^0(W, \text{ad}(E))$



$$W = I \times C$$

RESTRICT TO

$$C_y = y \times C$$

$$y \in I$$

ANY CONNECTION A

DETERMINES A HOLOMORPHIC BUNDLE

~~then~~

$$E_y = E|_{C_y}$$

$$E_y \rightarrow C_y$$

BOGOMOLNY EQNS ~~then~~

IMPLY THAT THE HOLOMORPHIC

TYPE OF E_y IS INDEPENDENT

OF y

$z =$ LOCAL

COORD ON C

$$F = d^2 z d\bar{z} F_{z\bar{z}} + \dots$$

$$D\phi = d\bar{z} \frac{D\phi}{D\bar{z}} + \dots$$

$$F_{y\bar{z}} = -i D_{\bar{z}} \phi$$

(101)

i.e.

$$\frac{\partial}{\partial y} (D_{\bar{z}}) = [D_{\bar{z}}, -i\phi]$$

THE $\bar{\partial}_A$ OPERATOR IS

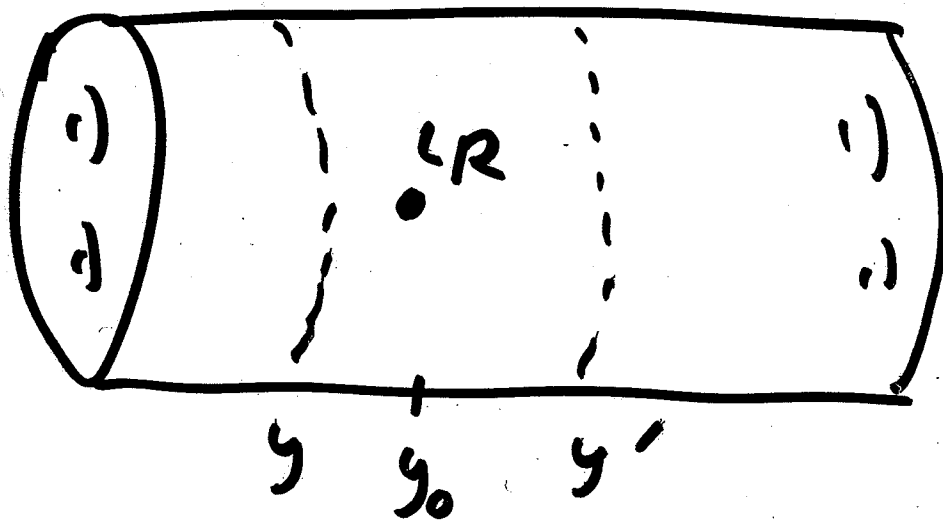
INDEPENDENT OF y

UP TO CONJUGACY.

$$\bar{\partial}_A = d\bar{z} D_{\bar{z}}$$

(102)

NOW IF AN 'E HOOFT OPERATOR
IS PRESENT AT $y_0 \times p \in W$



FOR $y_0 \in I$, $p \in \mathbb{C}$

THEN THE HOLOMORPHIC

TYPE OF E_y JUMPS AT

$y = y_0$.

$W \setminus (I \times p)$

BUT IT IS UNCHANGED

IF WE RESTRICT E_g

TO $C \setminus P$

THE CHANGE IN E_g IS

A "HEKE MODIFICATION"

OF TYPE $L_R \dots$

DETERMINED BY THE

$\frac{1}{2}$ HOOFT OPERATOR.

COISOTROPIC BRANE... KAP + ORLOR

~~GAN~~

ANY BRANE B_0

$\mathcal{A} = (B_0, B_0)$ STRINGS FORM A RING

\Rightarrow SHEAF OF RINGS

...

FOR ANY B

(B_0, B) STRINGS =

AN \mathcal{A} -MODULE

FIND "CANONICAL
COIS. BRANE B_0 "

set. $\mathcal{A} = \{ \text{SHEAF OF
DIFF OPS
ON } \mathcal{M} \}$

$T^*\mathcal{M} \subset \mathcal{M}^H$